# ON INTEGRABLE THREE-BODY PROBLEMS ON THE LINE 


#### Abstract

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The natural systems of three pair-interacting particles on the line are investigated. The properties of interactive potentials are considered under assumption that the given system has the first integral which is a polynomial of prescribed degree in the momenta. The functional equations for those potentials are obtained. All such potentials for special functional classes are described. А.Я. Вус. Интегрируемъе задачи трех тел на прямой// Математичні Студії. - 1998. - Т.10, № 1. - С.97-102.

Исследуются натуральные системы трех попарно взаимодействующих частиц на прямой. Рассмотрены свойства потенциалов взаимодействия в предположении, что данная система обладает первым интегралом, полиномиальным по импульсам. Получены функциональные уравнения на потенциалы и описаны все их решения для некоторых функциональных классов.


The dynamics of $n$ equal pair-interactive particles on the line is described by the Hamiltonian system with the Hamiltonian

$$
\begin{equation*}
H=\frac{1}{2} \sum_{i=1}^{n} p_{i}^{2}+\sum_{i<j} V\left(x_{i}-x_{j}\right) \tag{1}
\end{equation*}
$$

where the $x_{i}$ and $p_{i}, i=1, \ldots, n$, are the coordinates and momenta of the particles. We henceforth call the function $V$ a potential. We say that a potential $V$ admits an integral $F$ if $F$ is the first integral of the Hamiltonian system (1). We call the first integral $F$ to be nontrivial if $F$ is functionally independent with $H$. Such systems were considered in [1],[2] and the complete integrability for special cases of the Weierstrass $\wp$ function as the interaction potential was established. It is known ([3]), that the Hamiltonian system (1) is completely integrable for $V(x)$ being the Weierstrass $\wp$ function. A distinguishing feature of this problem is the polynomial character in the momenta of their additional integrals. It is therefore natural to obtain a description of Hamiltonians (1) which admits integrals that are polynomials in the momenta. In the paper this problem will be considered for $n=3$ and $V(x)$ satisfying tne following conditions:

1) $V$ is meromorphic in the vicinity of zero,
2) $V(x)=V(-x)$.

The total momentum $P=\sum p_{i}$ is the first integral of the Hamiltonian system (1). Therefore, this system can be reduced to the system with two degrees of freedom and the Hamiltonian

$$
\begin{equation*}
H=\frac{1}{2}\left(p_{1}^{2}+p_{2}^{2}\right)+W\left(x_{1}, x_{2}\right) \tag{2}
\end{equation*}
$$

where

$$
\begin{equation*}
W=V\left(x_{1}\right)+V\left(-\frac{x_{1}}{2}+\frac{x_{2} \sqrt{3}}{2}\right)+V\left(-\frac{x_{1}}{2}-\frac{x_{2} \sqrt{3}}{2}\right) . \tag{3}
\end{equation*}
$$

Without loss of generality we can assume that the first integral of system (2) has the form

$$
\begin{equation*}
F=F_{2 N}+F_{2 N-2}+\cdots+F_{0} \tag{4}
\end{equation*}
$$

$F_{k}=\sum_{i=1}^{k} E^{k-i, i}\left(x_{1}, x_{2}\right) p_{1}^{k-i} p_{2}^{i}$. Then the equation $\{F, H\}=0$ can be written in the form

$$
\begin{gather*}
0=p_{1} \partial_{x_{1}} F_{2 N}+p_{2} \partial_{x_{2}} F_{2 N}, \\
\partial_{p_{1}} F_{2 N} \cdot W_{1}+\partial p_{2} F_{2 N} \cdot W_{2}=p_{1} \partial_{x_{1}} F_{2 N-2}+p_{2} \partial_{x_{2}} F_{2 N-2},  \tag{5}\\
\ldots \ldots \ldots \ldots \ldots \\
\partial_{p_{1}} F_{2} \cdot W_{1}+\partial p_{2} F_{2} \cdot W_{2}=p_{1} \partial_{x_{1}} F_{0}+p_{2} \partial_{x_{2}} F_{0}
\end{gather*}
$$

where $\partial_{t}=\frac{\partial}{\partial t}, W_{i}=\frac{\partial W}{\partial x_{i}}$.
Let $R_{k}\left(x_{1}, x_{2}, p_{1}, p_{2}\right)=\sum f_{i}\left(x_{1}, x_{2}\right) p_{1}^{k-i} p_{2}^{i}$ be a homogeneous polynomial in the momenta. We shall put

$$
\left[R_{k}\right]=R_{k}\left(x_{1}, x_{2}, \partial_{x_{2}},-\partial_{x_{1}}\right)=\sum \partial_{x_{2}}^{k-i}\left(-\partial_{x_{1}}\right)^{i} f_{i}\left(x_{1}, x_{2}\right)
$$

Then equations (5) yield

$$
\begin{equation*}
P_{i}(V)=0, \tag{6}
\end{equation*}
$$

where $P_{i}(V)=\left[\partial_{p_{1}}\left(F_{2 N+2-2 i}\right) W_{1}\right]+\left[\partial_{p_{2}}\left(F_{2 N+2-2 i}\right) W_{2}\right], i=1, \ldots, N$.
Each nontrivial relation (6) is called the addition theorem.
The following result holds:
Lemma 1. The relation (6) is nontrivial for $i=1$ or $i=2$.
The proof of Lemma 1 is based on the consideration of relation (6) for $V(x)=$ $x^{-1}$. It is easy to show that if $\partial_{x_{1}} E^{0,2 N} \neq 0$, then $P_{1}(V) \neq 0$. Otherewise, $P_{2}(V) \neq 0$.

We shall consider only potentials satisfying the following conditions:
a) zero is a pole of order 2 for $V$;
b) $V$ is holomorphic in $\mathbb{R} \backslash\{0\}$;
c) $\lim _{x \rightarrow \infty} V(x)=0$.

The description of integrable potentials $V(x)$ in the class of functions under consideration is given in Theorems 1,2 .

Theorem 1. Let the leading homogeneous component of $F$ has constant coefficients. Then $V(x)$ is one of the Weierstrass $\wp$ function degenerated cases $x^{-2}$, $\sinh ^{-2} k x$.

The proof of Theorem 1 is based on three following lemmas.
Lemma 2. The addition theorem is $P_{2}(V)=0$ and $P_{2}(V)=[L] \varrho(V)$, where $L$ is a homogeneous polynomial in $p_{1}, p_{2}$ of degree $2 N-3$ and

$$
\begin{aligned}
\varrho(V) & =V\left(x_{1}\right) V^{\prime}\left(-\frac{x_{1}}{2}+\frac{x_{2} \sqrt{3}}{2}\right)-V^{\prime}\left(x_{1}\right) V\left(-\frac{x_{1}}{2}+\frac{x_{2} \sqrt{3}}{2}\right)+ \\
& +V\left(-\frac{x_{1}}{2}+\frac{x_{2} \sqrt{3}}{2}\right) V^{\prime}\left(-\frac{x_{1}}{2}-\frac{x_{2} \sqrt{3}}{2}\right)-V^{\prime}\left(-\frac{x_{1}}{2}+\frac{x_{2} \sqrt{3}}{2}\right) V\left(-\frac{x_{1}}{2}-\frac{x_{2} \sqrt{3}}{2}\right)+ \\
& +V\left(-\frac{x_{1}}{2}-\frac{x_{2} \sqrt{3}}{2}\right) V^{\prime}\left(x_{1}\right)-V^{\prime}\left(-\frac{x_{1}}{2}-\frac{x_{2} \sqrt{3}}{2}\right) V\left(x_{1}\right) .
\end{aligned}
$$

Proof. Consider the expression $F_{2 N}$. Without loss of generality we can assume that the integral $F$ is invariant with respect to the canonical transformation

$$
x_{i} \rightarrow \widetilde{x}_{i}, \quad p_{i} \rightarrow \widetilde{p}_{i}
$$

where

$$
\begin{array}{ll}
\widetilde{x}_{1}=-x_{1} / 2+x_{2} \sqrt{3} / 2, & \widetilde{x}_{2}=-x_{1} \sqrt{3} / 2-x_{2} / 2 \\
\widetilde{p}_{1}=-p_{1} / 2+p_{2} \sqrt{3} / 2, & \widetilde{p}_{2}=-p_{1} \sqrt{3} / 2-p_{2} / 2
\end{array}
$$

Thus we can represent $F_{2 N}$ in the form

$$
F_{2 N}=\left(p_{1} \widetilde{p}_{1} \widetilde{\widetilde{p}}_{1}\right)^{2} \cdot G(T, J)
$$

where $T=p_{1}^{2}+p_{2}^{2}, J=p_{2} \widetilde{p}_{2} \widetilde{\widetilde{p}}_{2}$ and $G$ is a polynomial in its variables. Then $P_{1}(V) \equiv 0$ and

$$
\begin{equation*}
F_{2 N-2}=\left(p_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial p_{1}} \cdot V\left(x_{1}\right)+\left(\widetilde{p}_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial \widetilde{p}_{1}} \cdot V\left(\widetilde{x}_{1}\right)+\left(\widetilde{\widetilde{p}}_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial \widetilde{p}_{1}} \cdot V\left(\widetilde{\widetilde{x}}_{1}\right)+\omega \tag{8}
\end{equation*}
$$

where $[\omega]=0$.
Then we get

$$
P_{2}(V)=\left[\left(3 J G-\left(p_{1} \widetilde{p}_{1} \widetilde{\widetilde{p}}_{1}\right)^{2} \frac{\partial G}{\partial J} p_{1} \widetilde{p}_{1} \widetilde{\widetilde{p}}_{1}\right] \varrho(V) .\right.
$$

The expression $L=\left(3 J G-\left(p_{1} \widetilde{p}_{1} \widetilde{\widetilde{p}}_{1}\right)^{2} \partial G / \partial J\right) p_{1} \widetilde{p}_{1} \widetilde{\widetilde{p}}_{1}$ is a homogeneous polynomial in $p_{1}, p_{2}$.

Lemma 3. (The uniqueness theorem). Let $[L] f=0$, where

$$
L=a p_{1}^{2}+b p_{1} p_{2}+c p_{2}^{2}, \quad a^{2}+b^{2}+c^{2} \neq 0
$$

$f$ is a holomorphic function in $\mathbb{R}^{2}$ and $\lim _{x^{2}+y^{2} \rightarrow \infty} f(x, y)=0$. Then $f \equiv 0$ in $\mathbb{R}^{2}$.
If $b^{2}-4 a c<0$, then this lemma is a generalization of the Liouville uniqueness theorem for the harmonic functions in $\mathbb{R}^{2}$.

If $b^{2}-4 a c \geq 0$, then $L$ can be decomposed into the product of two monomials. The equation $a \partial_{x} f+b \partial_{y} f=0$ yields $f(x, y)=g(b x-a y)$. Using the asymptotical condition we get $f \equiv 0$.

Since every homogeneous polynomial $L$ in $p_{1}, p_{2}$ can be represented as a product of polynomials of degree 1 or 2 , we receive from Lemmas 1,2 the equation $\varrho(V)=0$.

Lemma 4. The identity $\varrho(V)=0$ is the addition theorem for the integral of degree 3 and only the Weierstrass $\wp$ function satisfies it.
Proof. Consider the Laurent expansion at $x_{1}=0$ of the expression $\varrho(V)$. Its first nontrivial coefficient can be written in the form

$$
\begin{equation*}
V^{\prime \prime \prime}(t)-12 V(t) V^{\prime}(t)=0 \tag{9}
\end{equation*}
$$

where $t=x_{2} \sqrt{3} / 2$. It is known that the solution of equation (9) is either the Weierstrass $\wp$ function or one of its degenerated cases.

Theorem 2. Let the leading homogeneous component of $F$ has nonconstant coefficients and $\partial_{x_{1}} E^{0,2 N} \neq 0$. Then $V(x)=x^{-2}$.

Proof. It is easy to show that under conditions of Theorem 2 the relation $P_{1}(V)=0$ is nontrivial. Equations (5) yield

$$
\begin{aligned}
& F_{2 N-2}=\left(p_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial p_{1}} \cdot V\left(x_{1}\right)+\left(p_{1}\right)^{-2} \frac{\partial F_{2 N}}{\partial x_{1}} \cdot U\left(x_{1}\right)+\left(\widetilde{p}_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial \widetilde{p}_{1}} \cdot V\left(\widetilde{x}_{1}\right)+ \\
& \quad+\left(\widetilde{p}_{1}\right)^{-2} \frac{\partial F_{2 N}}{\partial \widetilde{x}_{1}} \cdot U\left(\widetilde{x}_{1}\right)+\left(\widetilde{\widetilde{p}}_{1}\right)^{-1} \frac{\partial F_{2 N}}{\partial \widetilde{\widetilde{p}}_{1}} \cdot V\left(\widetilde{\widetilde{x}}_{1}\right)+\left(\widetilde{\widetilde{p}}_{1}\right)^{-2} \frac{\partial F_{2 N}}{\partial \widetilde{\widetilde{x}}_{1}} \cdot U\left(\widetilde{\widetilde{x}}_{1}\right)+\omega
\end{aligned}
$$

where $V(t)=\partial_{t} U(t),[\omega]=0$.
Then the solution of the addition theorem $P_{1}=0$ can be written in the form

$$
\begin{equation*}
V(x)=\partial_{x}\left(Q_{2 N+4}(x) /\left(\partial_{x} E^{0,2 N}(x)\right)\right), \tag{10}
\end{equation*}
$$

where $Q_{2 N+4}$ is a polynomial of degree $2 N+4$ in $x$.
Lemma 5. Consider potential (10) as $V(z), z \in \mathbb{C}$. If $a \neq 0$ is a pole for $V(z)$, then $2 a$ is also a pole.

The proof of Lemma 5 is based on consideration of the Laurent expansion (at $x_{1}=a$ ) of the coefficients $E^{k-i, i}$. Then one can show that $z=a$ is a pole of order 2 for $V(z)$. Let

$$
B\left(x_{2}\right)=V\left(\frac{x_{2} \sqrt{3}}{2}+\frac{a}{2}\right)-V\left(\frac{x_{2} \sqrt{3}}{2}-\frac{a}{2}\right)
$$

Consider the order of the pole at $x_{1}=a$ of the coefficients $E^{2 N-4-i, i}$. The expression $B\left(x_{2}\right)$ can be represented as $B\left(x_{2}\right)=g\left(x_{2}\right)^{-3 /(2 m)}$, where $g\left(x_{2}\right)$ is a polynomial of degree $k \leq 2 m$.

Since $B\left(x_{2}\right)$ is a rational function, $B\left(x_{2}\right)=P_{3}^{-1}\left(x_{2}\right)$, where $P_{3}$ is a nontrivial polynomial of degree 3 . Let $z= \pm a$ be the poles and $z= \pm 2 a$ the regular points for $V(z)$. Then $B\left(x_{2}\right)$ has two poles of degree 3 at $x_{2}=a \sqrt{3}$. But the expression $P_{3}^{-1}\left(x_{2}\right)$ cannot have more then 3 poles in $\mathbb{C}$. This contradiction proves Lemma 5.

It follows from the previous lemma that potential (10) cannot have a pole $a \neq 0$. Therefore, $V(x)=x^{-2}+Q(x)$, where $Q(x)$ is a polynomial in $x$. In the class of potentials under consideration, $V(x)=x^{-2}$.

Let the potential $V$ satisfy only conditions (7a,7b). Then the following result is established:
Theorem 3. Let the potential $V(x)=x^{-2}+Q(x)$ admit integral (4). Then $\operatorname{deg} Q(x) \leq 4$.
Lemma 6. Let the potential $V(x)=x^{-2}+Q(x)$ admit an integral $F$ of degree $2 N$ and $\operatorname{deg} Q(x)=k$. Then the potential $x^{k}$ admits a nontrivial integral, which is polynomial of degree $d \leq 2 N$.

The proof of Lemma 6 is based on using of the canonical transformation

$$
x_{1}=u_{1} \cdot \varepsilon, x_{2}=u_{2} \cdot \varepsilon, \quad p_{1}=q_{1} \cdot \varepsilon^{-1}, p_{2}=q_{2} \cdot \varepsilon^{-1} .
$$

Then $H$ and $F$ are holomorphic functions of the parameter $\varepsilon$. Let $H_{0}=\left.H\right|_{\varepsilon=0}$, $F_{0}=\left.F\right|_{\varepsilon=0}$. Then $\left\{H_{0}, F_{0}\right\}=0$, where $H_{0}$ is a Hamiltonian for the system with $V(x)=x^{k}$ and $F_{0}$ is a nontrivial polynomial in $p_{1}, p_{2}$.

For proof of Theorem 3 we use Yoshida's theorem [4] on the nonintegrability of natural systems with the Hamiltonian (2), where $W$ is the homogeneous polynomial in $x_{1}, x_{2}$ of degree $k$.

Following Yoshida's method, consider the system

$$
\left\{\begin{array}{l}
\partial_{x_{1}} W=x_{1},  \tag{11}\\
\partial_{x_{2}} W=x_{2},
\end{array}\right.
$$

where

$$
W=x_{1}^{k}+\widetilde{x}_{1}^{k}+\widetilde{\widetilde{x}}_{1}^{k}
$$

Its solution is

$$
\left\{\begin{array}{l}
x_{1}=\left(2 /\left(k\left(1+2^{k-1}\right)\right)\right)^{1 /(k-2)} \\
x_{2}=x_{1} \sqrt{3}
\end{array}\right.
$$

For these $x_{1}, x_{2}$ the eigenvalues $\lambda_{1}, \lambda_{2}$ of the matrix $\left(\begin{array}{ll}W_{11} & W_{12} \\ W_{21} & W_{22}\end{array}\right)$ are

$$
\lambda_{1}=k-1, \quad \lambda_{2}=3(k-1) /\left(1+2^{k-1}\right)
$$

Calculating Kovalevska's indicators

$$
\varrho_{i}=\sqrt{1+8 k \lambda_{i} /(k-2)^{2}}
$$

we get $\varrho_{1}=(3 k-2) /(k-2) \in \mathbb{Q}$ and

$$
\varrho_{2}=\sqrt{1+\frac{24 k(k-1)}{(k-2)^{2}\left(1+2^{k-1}\right)}} .
$$

If $\varrho_{2} \notin \mathbb{Q}$, then it follows from Yoshida's theorem that the Hamiltonian system with $V(x)=x^{k}$ is nonintegrable.

Proposition. If $(k, m)$ is a solution of the Diophantine equation

$$
\begin{equation*}
(k-2)^{2}\left(1+2^{k-1}\right)^{2}+24 k(k-1)\left(1+2^{k-1}\right)=m^{2}, \tag{12}
\end{equation*}
$$

then $k \leq 4$.
Proof. For $k=6,8$ or 10 we can directly calculate $m$. Consider the case $k>10$. Let $(k, m)$ be a solution of (12). One can show that

$$
m=2(l-1)\left(1+2^{2 l-1}\right)+12 s
$$

where $k=2 l$ and $s$ is an integer number.Then $s$ satisfies the Diophantine equation

$$
(l(2 l-1)-s(l-1))\left(1+2^{2 l-1}\right)=3 s^{2} .
$$

But for $l \geq 6$ its solution must satisfy the inequalities $2 l+1<s<2 l+2$.
The last proposition completes the proof of Theorem 3.

## REFERENCES

[1] Moser J., Three integrable Hamiltonian systems connected with semisimple Lie algebras, Invent. Math. 37 (1976), no. 2, 93-108.
[2] Calogero F, Exactly solvable one-dimensional many body problems, Letters al Nuovo Cimento 13 (1975), no. 11, 411-416.
[3] Pidkuiko S.I., Stepin A.M., Polynomial integrals of Hamiltonian systems, Dokl. Akad. Nauk SSSR 239 (1978), no. 1, 50-51, English transl. in Soviet Math. Dokl. 19 (1978), no.2, 282-286.
[4] Yoshida H., A criterion for the non-existence of an additional analytic integral in Hamiltonian systems with $n$ degrees of freedom, Phys.Lett.A. 141 (1989), no. 3-4, 108-112.

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